

The field content of the Minimal Supersymmetric Standard Model (MSSM)

by M.B.Kocic – Version 1.05 (2016-01-25) – SUSYRC-HT15

	Vector supermultiplets	Superfield	Adj. repr.	Spin-1 (gauge bosons)	Spin-1/2 (gauginos)	Aux.	Vector superfield (in Wess-Zumino gauge)
Gauge fields	$U(1)_Y$	V_Y	$\mathbf{1}, \mathbf{1}, 0$	B_μ , B-boson	$\lambda_Y \equiv \tilde{B}$, bino	D_Y	$V_Y \equiv \theta \sigma^\mu \bar{\theta} B_\mu + \theta\theta \bar{\theta} \tilde{\lambda}_Y + \bar{\theta}\bar{\theta} \theta \lambda_Y + \frac{1}{2} \theta\theta \bar{\theta}\bar{\theta} D_Y$
	$SU(2)_L$	V_L^i	$\mathbf{1}, \mathbf{3}, 0$	W_μ^i , W-bosons	$\lambda_L^i \equiv \tilde{W}^i$, winos	D_L^i	$V_L^i \equiv \theta \sigma^\mu \bar{\theta} W_\mu^i + \theta\theta \bar{\theta} \tilde{\lambda}_L^i + \bar{\theta}\bar{\theta} \theta \lambda_L^i + \frac{1}{2} \theta\theta \bar{\theta}\bar{\theta} D_L^i$
	$SU(3)_C$	V_C^a	$\mathbf{8}, \mathbf{1}, 0$	g_μ^a , gluons	$\lambda_C^a \equiv \tilde{g}^a$, gluinos	D_C^a	$V_C^a \equiv \theta \sigma^\mu \bar{\theta} g_\mu^a + \theta\theta \bar{\theta} \tilde{\lambda}_C^a + \bar{\theta}\bar{\theta} \theta \lambda_C^a + \frac{1}{2} \theta\theta \bar{\theta}\bar{\theta} D_C^a$
	Chiral supermultiplets	Superfield	Repr.	Spin-1/2 (fermions)	Spin-0 (sfermions)	Aux.	Chiral superfield (in terms of $y^\mu = x^\mu - i\theta\sigma^\mu\bar{\theta}$)
Matter fields	quarks, s(calar) quarks	Q_I	$\mathbf{3}, \mathbf{2}, +\frac{1}{6}$	$\begin{pmatrix} u_L \\ d_L \end{pmatrix}, \begin{pmatrix} \chi_u \\ \chi_d \end{pmatrix}$	$\begin{pmatrix} \tilde{u}_L \\ \tilde{d}_L \end{pmatrix}, \begin{pmatrix} \tilde{\phi}_u \\ \tilde{\phi}_d \end{pmatrix}$	$\begin{pmatrix} F_u \\ F_d \end{pmatrix}$	$Q = Q_1 = \begin{pmatrix} \tilde{\phi}_u + \sqrt{2}\theta\chi_u + \theta\theta F_u \\ \tilde{\phi}_d + \sqrt{2}\theta\chi_d + \theta\theta F_d \end{pmatrix}$
		$U_I (\bar{u}_I)$	$\bar{\mathbf{3}}, \mathbf{1}, -\frac{2}{3}$	$\bar{u}_L = (u_R)^c, \chi_{\bar{u}}$	$\tilde{u}_L, \tilde{\phi}_{\bar{u}}$	$F_{\bar{u}}$	$\bar{u} = U_1 = \tilde{\phi}_{\bar{u}} + \sqrt{2}\theta\chi_{\bar{u}} + \theta\theta F_{\bar{u}}$
		$D_I (\bar{d}_I)$	$\bar{\mathbf{3}}, \mathbf{1}, +\frac{1}{3}$	$\bar{d}_L = (d_R)^c, \chi_{\bar{d}}$	$\tilde{d}_L, \tilde{\phi}_{\bar{d}}$	$F_{\bar{d}}$	$\bar{d} = D_1 = \tilde{\phi}_{\bar{d}} + \sqrt{2}\theta\chi_{\bar{d}} + \theta\theta F_{\bar{d}}$
	leptons, s(calar) leptons	L_I	$\mathbf{1}, \mathbf{2}, -\frac{1}{2}$	$\begin{pmatrix} \nu_{eL} \\ e_L \end{pmatrix}, \begin{pmatrix} \chi_{\nu_e} \\ \chi_e \end{pmatrix}$	$\begin{pmatrix} \tilde{\nu}_{eL} \\ \tilde{e}_L \end{pmatrix}, \begin{pmatrix} \tilde{\phi}_{\nu_e} \\ \tilde{\phi}_e \end{pmatrix}$	$\begin{pmatrix} F_{\nu_e} \\ F_e \end{pmatrix}$	$L = L_1 = \begin{pmatrix} \tilde{\phi}_{\nu_e} + \sqrt{2}\theta\chi_{\nu_e} + \theta\theta F_{\nu_e} \\ \tilde{\phi}_e + \sqrt{2}\theta\chi_e + \theta\theta F_e \end{pmatrix}$
		$N_I (\bar{\nu}_I)$	$\mathbf{1}, \mathbf{1}, 0$	$\bar{\nu}_{eL} = (\nu_{eR})^c, \chi_{\bar{\nu}_e}$	$\tilde{\nu}_{eL}, \tilde{\phi}_{\bar{\nu}_e}$	$F_{\bar{\nu}_e}$	$\bar{\nu} = N_1 = \tilde{\phi}_{\bar{\nu}_e} + \sqrt{2}\theta\chi_{\bar{\nu}_e} + \theta\theta F_{\bar{\nu}_e}$
		$E_I (\bar{e}_I)$	$\mathbf{1}, \mathbf{1}, +1$	$\bar{e}_L = (e_R)^c, \chi_{\bar{e}}$	$\tilde{e}_L, \tilde{\phi}_{\bar{e}}$	$F_{\bar{e}}$	$\bar{e} = E_1 = \tilde{\phi}_{\bar{e}} + \sqrt{2}\theta\chi_{\bar{e}} + \theta\theta F_{\bar{e}}$
Higgs fields	higgsinos, higgs	H_u	$\mathbf{1}, \mathbf{2}, +\frac{1}{2}$	$\begin{pmatrix} \tilde{H}_u^+ \\ \tilde{H}_u^0 \end{pmatrix}$	$\begin{pmatrix} H_u^+ \\ H_u^0 \end{pmatrix}$	$\begin{pmatrix} F_{H_u}^+ \\ F_{H_u}^0 \end{pmatrix}$	$H_u = \begin{pmatrix} H_u^+ + \sqrt{2}\theta\tilde{H}_u^+ + \theta\theta F_{H_u}^+ \\ H_u^0 + \sqrt{2}\theta\tilde{H}_u^0 + \theta\theta F_{H_u}^0 \end{pmatrix}$
		H_d	$\mathbf{1}, \mathbf{2}, -\frac{1}{2}$	$\begin{pmatrix} \tilde{H}_d^0 \\ \tilde{H}_d^- \end{pmatrix}$	$\begin{pmatrix} H_d^0 \\ H_d^- \end{pmatrix}$	$\begin{pmatrix} F_{H_d}^+ \\ F_{H_d}^0 \end{pmatrix}$	$H_d = \begin{pmatrix} H_d^0 + \sqrt{2}\theta\tilde{H}_d^0 + \theta\theta F_{H_d}^+ \\ H_d^- + \sqrt{2}\theta\tilde{H}_d^- + \theta\theta F_{H_d}^0 \end{pmatrix}$

N.B. Quantum numbers of the fields in the MSSM are the same as in the SM.

Superspace Formalism

A **vector superfield**, $SU(n)$, in WZ gauge:

$$V^i(x, \theta, \bar{\theta}) \equiv \theta\sigma^\mu\bar{\theta} V_\mu^i(x) + \theta\theta\bar{\theta}\tilde{\lambda}^i(x) + \bar{\theta}\bar{\theta}\theta\lambda^i(x) + \frac{1}{2}\theta\theta\bar{\theta}\bar{\theta} D^i(x), \quad V \equiv V^i T^i$$

The **field strength superfield**:

$$\mathcal{V}_\alpha \equiv \mathcal{V}_\alpha^i T^i \equiv -\frac{1}{4} \bar{D}_{\dot{\beta}} \bar{D}^{\dot{\beta}} (e^{-2gV} D_\alpha e^{2gV})$$

The gauge kinetic term:

$$\frac{1}{4} (\mathcal{V}^{i\alpha} \mathcal{V}_\alpha^i)|_F + \text{h.c.}$$

$\frac{1}{4} (\mathcal{V}^{i\alpha} \mathcal{V}_\alpha^i)|_F$ (mod tot. der.) =

$$\frac{1}{2} D^i D^i + i \bar{\lambda}^i \bar{\sigma}^\mu D_\mu \lambda^i - \frac{1}{4} F_{\mu\nu}^i F^{i\mu\nu}$$

The chiral covariant derivatives:

$$\begin{aligned} D_\alpha &\equiv \partial_\alpha - i\sigma_{\alpha\dot{\alpha}}^\mu \bar{\theta}^{\dot{\alpha}} \partial_\mu, \\ D^\alpha &\equiv -\partial^\alpha + i\bar{\theta}_{\dot{\alpha}} \bar{\sigma}^{\mu\dot{\alpha}\alpha} \partial_\mu, \\ \bar{D}_{\dot{\alpha}} &\equiv \bar{\partial}_{\dot{\alpha}} - i\bar{\theta}_{\dot{\alpha}} \bar{\sigma}^{\mu\dot{\alpha}\alpha} \partial_\mu, \\ \bar{D}_{\dot{\alpha}} &\equiv -\bar{\partial}_{\dot{\alpha}} + i\theta^\alpha \sigma_{\alpha\dot{\alpha}}^\mu \partial_\mu \end{aligned}$$

Observe:

$$\begin{aligned} \partial^\alpha &= -\varepsilon^{\alpha\beta} \partial_\beta \\ \bar{\partial}_{\dot{\alpha}} &= -\varepsilon^{\dot{\alpha}\dot{\beta}} \bar{\partial}_{\dot{\beta}} \end{aligned}$$

A **chiral superfield**:

$$\Phi(y, \theta) = \phi(y) + \sqrt{2}\theta\chi(y) + \theta\theta F(y)$$

where $y^\mu = x^\mu - i\theta\sigma^\mu\bar{\theta}$. (*) Expanded:

$$\begin{aligned} \Phi(x, \theta, \bar{\theta}) &= \phi(x) + \sqrt{2}\theta\chi(x) + \theta\theta F(x) \\ &\quad - \frac{1}{4}\theta\theta\bar{\theta}\bar{\theta}\square\phi(x) - \frac{i}{\sqrt{2}}\theta\theta\bar{\theta}\sigma^\mu\partial_\mu\chi(x) \end{aligned}$$

The canonical and gauge inv. kinetic terms:

$$\text{free: } (\Phi^\dagger\Phi)|_D, \quad \text{matter: } (\Phi^\dagger e^{2gV}\Phi)|_D$$

$(\Phi^\dagger\Phi)|_D$ (mod tot. der.) =

$$\partial_\mu\phi^\dagger\partial^\mu\phi + i\bar{\chi}\bar{\sigma}^\mu\partial_\mu\chi + F^\dagger F$$

$U(1)$: $(\Phi^\dagger e^{2gV}\Phi)|_D$ (mod tot. der.) =

$$\begin{aligned} &(D_\mu\phi)^\dagger D^\mu\phi + i\bar{\chi}\bar{\sigma}^\mu D_\mu\chi + F^\dagger F \\ &- g\phi^\dagger\phi D - (\sqrt{2}g\chi\lambda\phi^\dagger + \text{h.c.}) \end{aligned}$$

(*) y are complex composite bosonic coordinates.

The **superpotential** for the collection $\{\Phi_I\}$:

$$\mathcal{W}(\Phi_I) = \frac{1}{2} m^{IJ} \Phi_I \Phi_J + \frac{1}{6} y^{IJK} \Phi_I \Phi_J \Phi_K$$

Opt. F -term (as $c^I \Phi_I$ in the superpotential):

$$-k\Phi|_F = -kF \quad (\text{not in MSSM})$$

The interaction term:

$$\mathcal{W}(\Phi_I)|_F + \text{h.c.}$$

$\mathcal{W}(\Phi_I)|_F + \text{h.c.} =$

$$\frac{\partial}{\partial\phi_I} \mathcal{W}(\phi_I) F_I - \frac{1}{2} \frac{\partial^2}{\partial\phi_I \partial\phi_J} \mathcal{W}(\phi_I) \chi_I \chi_J + \text{h.c.}$$

Opt. Fayet-Iliopoulos D -term (in $U(1)$ only)

$$-2\kappa V|_D = -\kappa D \quad (\text{not in MSSM})$$

The MSSM Lagrangian

The (matter) kinetic term

$$\mathcal{L}^{\text{kin}} = \sum_{\Phi, gV} \Phi^\dagger e^{2gV} \Phi|_D$$

gV runs over all the gauge superfields,

$$\{gV_Y, g_i V_L^{i\dot{t}}, g_a V_C^a T^a\},$$

Φ runs over all the matter superfields,

$$\{Q_I, U_I, D_I, L_I, N_I, E_I, H_u, H_d\},$$

in the appropriate representation where $I = 1, 2, 3$ is the family index.

The gauge kinetic term

$$\mathcal{L}^{\text{g-kin}} = \sum_V \frac{1}{4} (\mathcal{V}^\alpha \mathcal{V}_\alpha)|_F + \text{h.c.}$$

\mathcal{V} runs over all the field strength superfields

The Yukawa interaction term

$$\mathcal{L}^{\text{Yuk}} = \mathcal{W}(\Phi)|_F + \text{h.c.}$$

\mathcal{W} is the superpotential (given below)

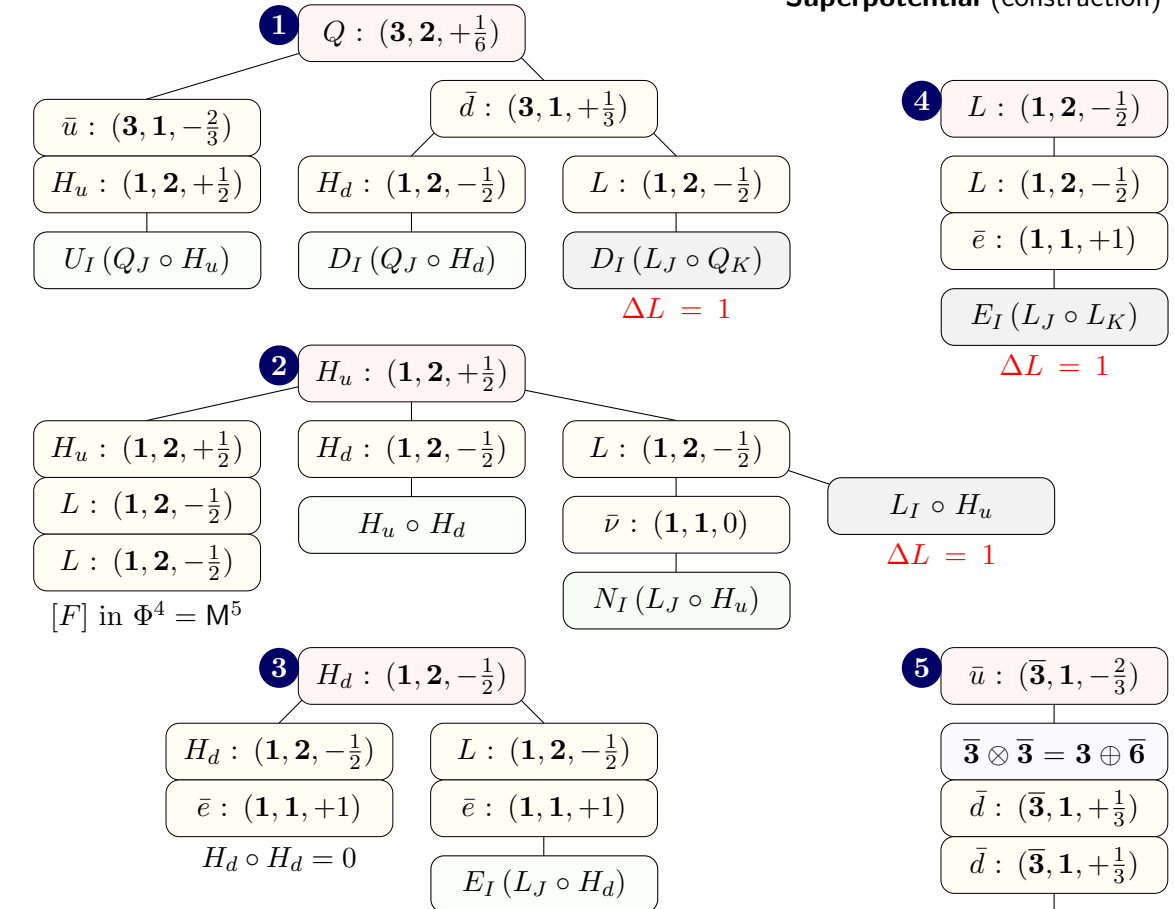
The Superpotential, Part I ('good' terms)

$$\begin{aligned} \mathcal{W}_1 &= y_u^{IJ} U_I (Q_J \circ H_u) - y_d^{IJ} D_I (Q_J \circ H_d) \\ &\quad + y_\nu^{IJ} N_I (L_J \circ H_u) - y_e^{IJ} E_I (L_J \circ H_d) \\ &\quad + \mu (H_u \circ H_d) \end{aligned}$$

The Superpotential, Part II (terms that violate lepton/baryon numbers)

$$\begin{aligned} \mathcal{W}_2 &= \lambda^{JK} E_I (L_J \circ L_K) + \lambda' IJK D_I (L_J \circ Q_K) \\ &\quad + \mu_0^I (L_I \circ H_u) + \lambda_3^IJK U_I D_J D_K \end{aligned}$$

Superpotential (construction)



$\circ \equiv SU(2)$ invariant product of doublets:

$$\begin{pmatrix} \Phi_1 \\ \Phi_2 \end{pmatrix} \circ \begin{pmatrix} \Psi_1 \\ \Psi_2 \end{pmatrix} \equiv \Phi_1 \Psi_2 - \Phi_2 \Psi_1$$

Useful formulae:

$$\mathcal{F}(\Phi)|_D \equiv \int d^2\theta d^2\bar{\theta} \mathcal{F}(\Phi),$$

$$\mathcal{F}(\Phi)|_F \equiv \int d^2\theta \mathcal{F}(\Phi),$$

In the WZ gauge: $e^V = 1 + V + \frac{1}{2}V^2$,

$$V^2 = \frac{1}{2}\theta\theta\bar{\theta}\bar{\theta} V_\mu V^\mu, \quad V^n = 0 \quad (n \geq 3).$$